

# Damping of Josephson oscillations in strongly correlated one-dimensional atomic gases

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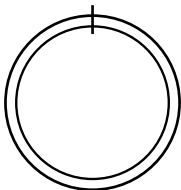
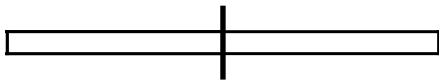


## Frank Hekking Memorial Workshop

January 28th, 2018



## Coupled finite wires and ring



### Wires

Infinite 1D coupled wires have been previously investigated showing dissipation [1]

[1] C. L. Kane, M. P. A. Fisher, *Phys. Rev. B* **46** 15233 (1992); G. Schön, A. D. Zaikin, *Phys. Reports* **198**, 237-412 (1990)

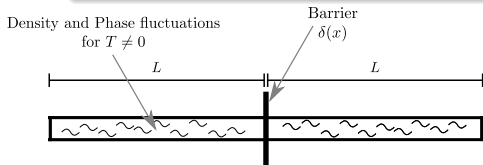
### Rings

Closed loops present new possibilities, especially for studying current dynamics and superfluidity [2]

[2] M. Cominotti, *et. al.*, *Phys. Rev. Lett.* **113**, 025301 (2014); D. Aghamalyan, *et. al.*, *New J. Phys.* **17** 045023 (2015)



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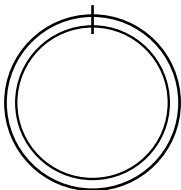
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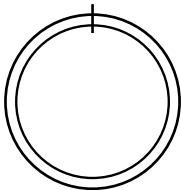
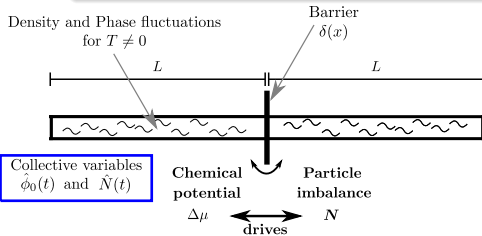
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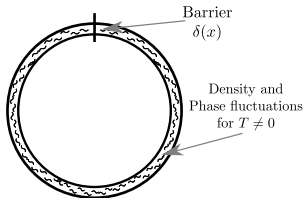
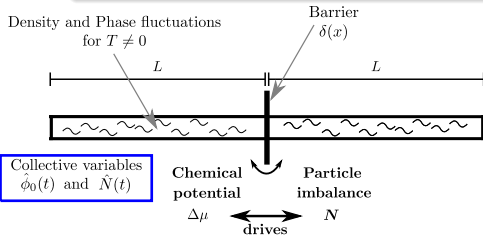
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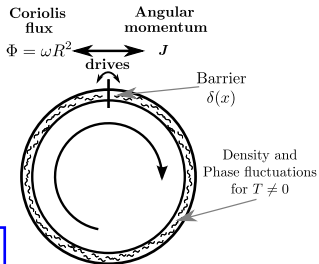
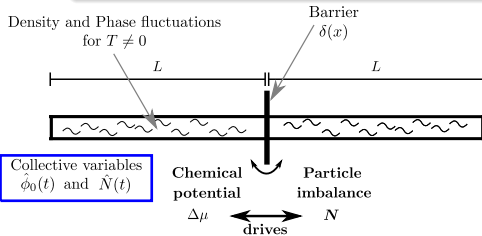
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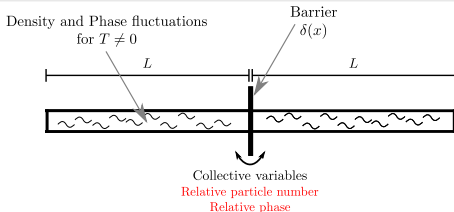
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## 1D Finite wires with strong barrier



### Luttinger Hamiltonian

$$H_{LL\pm} = \frac{\hbar v K}{2} \int_0^{\pm L} dx \left[ (\partial_x \phi_{\pm}(x, t))^2 + \frac{1}{K^2} (\partial_x \theta_{\pm}(x, t))^2 \right]$$
$$H_b = E_J \cos(\phi_+(0) - \phi_-(0))$$

## 1D Finite wires with strong barrier

Density and Phase fluctuations for  $T \neq 0$

Barrier  $\delta(x)$

Collective variables  
Relative particle number  
Relative phase

**Effective Hamiltonian**

Quantum particle

Bath

Interaction

$$\hat{H} = \underbrace{\frac{\hbar^2}{2ML^2} \hat{N}^2}_{\text{"Potential term"}} - \underbrace{E_J \cos(\hat{\phi}_0)}_{\text{"Nonlinear Kinetic term"}} + \sum_{\mu \geq 1} \left[ \underbrace{\frac{\hat{P}_\mu^2}{2M} + \frac{1}{2} M \Omega_\mu^2 \hat{Q}_\mu^2}_{\text{"Harmonic bath"}} + \underbrace{\frac{\sqrt{2}\hbar}{ML} \hat{N} \hat{P}_\mu}_{\text{"Coupling"}} + \underbrace{\frac{\hbar^2}{ML^2} \hat{N}^2}_{\text{"Renormalization"}} \right]$$

$$M = \frac{\hbar K}{2\pi vL} = \frac{K^2}{2\pi^2 N_0} m \quad \Omega_\mu = vk_\mu$$

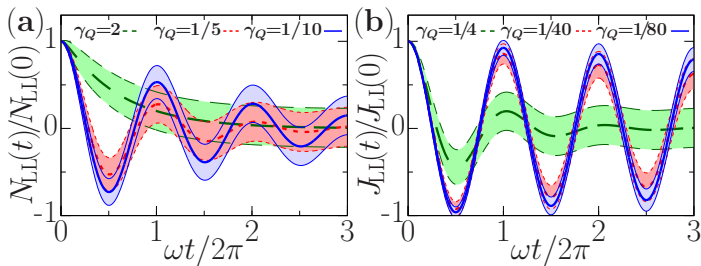




## Classical limit results within the Luttinger liquid approach

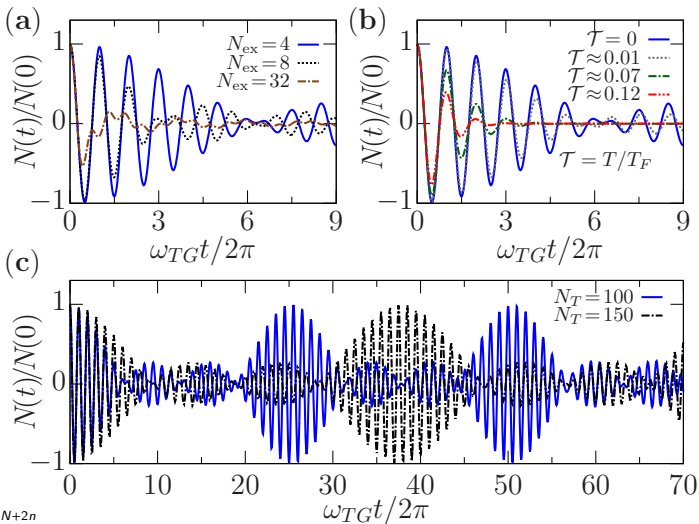
Quantum particle coupled to a bath  $\implies$  Quantum Langevin equation of motion

Classical limit shows the main dynamical features





## Exact Tonks-Girardeau method



$$\omega_{TG} = \frac{1}{N_{\text{ex}}} \sum_{n=-N_{\text{ex}}/2+1}^{N_{\text{ex}}/2} \omega_{N+2n}$$